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Thermal-hydraulic performance of a slit fin and influence of the fin pitch

Comportamiento termo hidráulico de una aleta tipo slit e influencia del paso entre aletas

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Abstract: The airside performance of compact heat exchanger with slit fins having one row of tubes is investigated in this study using computational fluid dynamics. A numerical approach was implemented and it was certified using correlations of published literature. The objective was to obtain the thermal-hydraulic behavior for different fin pitches while ranging the Reynolds number from 400 to 1800. The influence on heat transfer coefficient, fin efficiency and heat flux was found. It is observed to be a function of the fin pitch, which is at its best when the offset strips of one fin have the same distance to the ones of the opposite arranged strips as to that fin plate. Was found, for each fin pitch, a peak in the fin efficiency as a function of the inlet velocity.

Keywords: slit fin, compact heat exchanger, fin pitch, thermal-hydraulic behavior.

Resumen: El trabajo constituyó una aproximación numérica del desempeño térmico de intercambiadores de calor compactos con aletas slit y una sola fila de tubos. El modelo numérico se validó frente a correlaciones experimentales publicadas en la literatura. El objetivo fue estudiar el comportamiento termo-hidráulico como una función del espaciamiento entre aletas en un rango de números de Reynolds variable entre 400 y 1800. Se determinó la influencia del paso entre aletas sobre el coeficiente de transferencia de calor, la eficiencia de la aleta y el flujo de calor. Se observó que el paso entre aletas ejerce influencia sobre estas magnitudes encontrándose que el mejor valor era aquel que permitía que la parte horizontal de las slit quedase en el centro del canal entre aletas y a igual distancia de cada aleta. Además, se encontró que la eficiencia de la aleta posee un pico para cada paso entre aletas que depende de la velocidad de entrada.

Palabras claves: aleta slit, intercambiador de calor, paso entre aletas, comportamiento termo hidráulico.

Introduction

A high-energy loss in most industrial or domestic processes is caused in the heat transfer process due to the temperature difference between the fluids. Often one of the working fluids is air since it is a clean, cheap and stable material but usually comes with a high thermal resistance. To reduce fuel and power consumption as well as to lower climate-damaging gas emissions compact heat exchangers regularly use interrupted surfaces. Wavy, louver, slit fins and many more are widely used to increase the thermal performance by preventing the formation of thick boundary

layers and creating high turbulence in the heat exchanger. The effects of fin pitch, number of tube row and other geometrical characteristics have been investigated extensively and still there are researches going on. A lot of them are of experimental nature but since computing power is increasing also numerical studies using computational fluid dynamics (CFD) are elaborated like the current article.

Wang *et al.* [1] conducted an experimental study on the air side performance of compact slit fin-and-tube heat exchangers and found that it is relatively independent on the number of tube rows and the fin pitch. However, they only studied two different fin pitches in this work. Using various comparison methods, they evaluated the results with the louver and plain fin. They discovered that the air side performance of interrupted fins is superior to that of the plain fin but this difference decreases with lower fin pitch and Reynolds number. An updated correlation was proposed and is considered in the current paper to certify the numerical model with experimental data.

However in a different article Wang *et al.* [2] studied experimentally the airside performance of fin-and-tube heat exchangers, having a different slit geometry and came to the conclusion that the heat transfer performance increases with decrease of fin pitch for only one tube row (N). For $N > 4$, the effect of fin pitch on the heat transfer performance was reversed and the heat transfer performance decreased with increased number of tube row. The friction factor was found relatively independent to the number of tube row. Borrajo-Perez *et al.* [3] investigated the thermal-hydraulic characterization of a compact heat exchanger having two rows of tubes and rectangular wavy fins and used the proposed correlations in Wang *et al.* [1] to certify the numerical procedure and its mathematical model against published literature.

Also Wang *et al.* [4] carried out a study on convex louver fins where the heat transfer data was presented as a multiplication of the fin efficacy η_f by the average heat transfer coefficient $\bar{h}$. The separation of the fin efficiency and the heat transfer coefficient were based on the Schmidt method, which also will be used in this paper. The Colburn and Fanning friction factor are calculated using the correlations utilized by Wang *et al.* [5] in an experimental study on the airside performance of a herringbone wavy fin. The conclusions for modeling and simulating a compact heat exchanger will be used in the current paper. Interrupted fins have a quite complex geometry and to study the heat transfer coefficient the fin efficiency is not determined precisely by the Schmidt method, which performs better for continuous fin designs. This is because the heat transfer conduction inside the solid is disturbed and several changes in the fin temperature distribution can be found in interrupted fins. In numerical studies, it is possible to determine the fin efficiency and heat transfer coefficient directly. Unfortunately many papers working with interrupted fins still are using the approximate Schmidt method [6]

In the study of Ameel *et al.* [7] was shown that the Schmidt correlation for a fin having a complex, interrupted design results in an over estimation of the fin efficiency. They proposed a new numerical method other than

From the literature reviewed can be noted that is not clear the influence of the fin pitch on the thermo-hydraulic behavior of slit fin having one row of tubes. On the other hand, the fin efficiency using the Schmidt method is not appropriate for interrupted fins. The objective was to obtain the thermal-hydraulic behavior for different fin pitches while ranging the Reynolds number from 400 to 1800 using a novel reported method for fin efficiency determination [7].

In this study, a model of a fin-and-tube heat exchanger having slit fin geometry is investigated. The related geometry is illustrated in figure 1. The present fin has an offset slit geometry and for an accurate differentiation, the plate fin will be just called “fin” and the offsets “strips”. After the numerical validation, a model with slightly bigger strips and three different fin pitches is examined.



Computational Domain

The geometry of the fin used in this work was modeled in the Inventor software. A model having only one row of tubes is presented in this work. For the three-dimensional physical domain, a fin was set in the middle and

consecutive one-half of the fin pitch in both directions. When the heat transfer and the flow conditions are periodic in the upper and lower side of the fin, it is very simple to implement this boundary condition in the two parallel planes at the top and the bottom of the flow domain. The size and the design of the model are similar as in the real heat exchanger. In width, there was only figure. 2 presents the computational domain with the boundary conditions labeled. The fin can be observed in the middle of the channel and its surface is shaded.

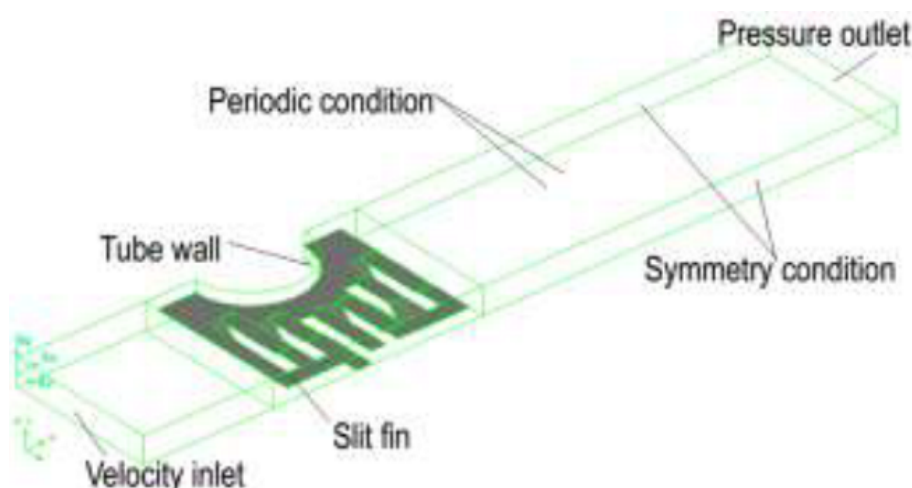


Fig.2
Computational domain in Gambit with the boundary conditions labeled.

Two regions were artificially extended in the computational domain. The one at the up stream is as long as the channels width for considering the inlet effects and downstream of the fins trailing edge the channel was extended three times its width. The extensions of the computational domain have a big impact on computational time but are necessary to show exit effects and avoid reversal flow at the exit. Two regions with symmetrical boundary conditions were established at both sides of the domain and two periodic boundary condition were set at top and bottom. A velocity inlet condition at inlet section and pressure outlet condition at outlet section were established. A boundary condition of constant temperature at 286K was used at the tube wall.

On the fins surface a wall-coupled boundary condition was applied to consider the conjugate heat transfer on the surface. The computational domain was meshed in the GAMBIT software and the grid quality was checked. Several sub domains were used to make the grid generation smoother. Approximately $5,4 \times 10^5$ tetrahedral volumes were generated for the finest mesh with 0,115mm length of the edges which is also the fin thickness (ft). Heat transfer and fluid flow simulations were performed using the commercial solver FLUENT. The material properties are found in table 1 and are considered constant.

Table 1
Material properties used in the numerical simulation

Material	Density (kg/m ³)	Specific heat capacity (J/kgK)	Thermal conductivity (W/mK)	Viscosity (N-s/m ²)
Air	1,225	1006,43	0,0242	1,789E-5
aluminum	2719	871	202,4	-

The air is assumed incompressible and the viscous model laminar. During the calculation, a segregated solver with constant properties was considered and to ensure mass conservation as well as to obtain a pressure field the SIMPLE algorithm for the coupling between pressure and velocity was implemented. Scale residuals were monitored and convergence criteria was set to 10^{-8} for energy and 10^{-4} for continuity. The discretization applied was standard for pressure, first order upwind for momentum and second order upwind for energy. The second order scheme was selected considering the misalignment between the flow and the grid. This decision helps to avoid the effect of numerical diffusion in misaligned grids. In general, the time elapsed for every solution was less than one hour in a specialized computer.

Data Reduction Procedure

The Reynolds number (Re) is determined using the velocity u_{max} at the minimal area of the channel. The characteristic length is the fin collar diameter D_c . Heat transferred in a heat exchanger ($\dot{Q}$) can be calculated with the change in temperature of the air between the inlet and outlet sections ($T_{out} - T_{in}$), the mass flow through the heat exchanger channel ($\dot{m}$) and knowing the fluid's specific heat capacity c_p . Equation 1

$$\dot{Q} = \dot{m} c_p (T_{out} - T_{in}) \quad (1)$$

When a fluid experiences a phase change inside tubes it is a common practice to consider an elevated value for the film heat transfer coefficient inside the tubes. This assumption is important because a global heat transfer coefficient can be obtained considering only the external film heat transfer coefficient and perfect conduction inside the tube wall. The wall temperature is set constant and with the same value than the refrigerant temperature flowing inside the tubes.

$$\Delta T_{log} = \frac{\Delta T_1 - \Delta T_2}{\ln \frac{\Delta T_1}{\Delta T_2}} \quad (2)$$

The LMTD correction factor F is considered unitary because one of the fluids flowing through the heat exchanger keeps a constant temperature.

The global heat transfer coefficient is calculated considering the equality of the heat calculated by the equations 1 and 3. The efficacy η_0 of the fin is involved and at the same time, it is a function of the global heat transfer coefficient. For this, the efficacy can be calculated if the fin efficiency η is known. A_f/A_0 is the rate of heat transfer area to the total area of the fin, as in equation 4.

The heat transferred (equation 4) can be calculated using the logarithmic mean temperature difference LMTD in equation 2 and by knowing the heat transfer area A_f and the global heat transfer coefficient.

$$\dot{Q}_0 = \eta_0 \bar{h} A_f F \Delta T_{\log} \quad (3)$$

$$\eta_0 = 1 - \frac{A_f}{A_0} (1 - \eta) \quad (4)$$

The fin efficiency for the rectangular fin is determined using an approximate method developed by Schmidt for circular fins and further explained by Wang *et al.* [4]. The fin efficiency is expressed by equation 5:

$$\eta = \frac{\tanh(m r_t \varphi)}{m r_t \varphi} \quad (5)$$

the parameter m is calculated with the thermal conductivity (k_f) of the fin and the fin thickness (f_t) equation 6:

$$m = \sqrt{\frac{2\bar{h}}{k_f f_t}} \quad (6)$$

The term φ is obtained with the equivalent tube radius divided by the tube radius (R_{eq}/r_t). This parameter is just depending on the heat exchanger geometry. Equation 7

$$\varphi = \left(\frac{R_{eq}}{r_t} - 1 \right) \left(1 + 0.35 \ln \left(\frac{R_{eq}}{r_t} \right) \right) \quad (7)$$

With, equation 8:

$$\frac{R_{eq}}{r_t} = 1.28 \frac{X_M}{r_t} \left(\frac{X_L}{X_M} - 0.2 \right)^{0.5} \quad (8)$$

Where X_M is the half-transversal pitch and X_L is calculated by the following equation, where S_T and S_L are the transversal and longitudinal pitch respectively equation 9:

$$X_L = (1/2) \sqrt{\left(\frac{S_T}{2}\right)^2 + (S_L)^2} \quad (9)$$

The fin efficacy and the global heat transfer coefficient have an implicit formulation; there for, an iterative procedure is needed. The condition of equality of the heats in equation 5 must be achieved. For a fixed geometry, there is only one pair of values in fin efficacy and global heat transfer coefficient, which meets this requirement. The Colburn factor j can be calculated having the global heat transfer coefficient, the velocity in the minimal cross section of the channel u_{max} , the heat capacity at constant pressure of the air, the Prandtl number (Pr) and the average air density equation 10:

$$j = \frac{\bar{h}}{\rho_m C_{pm} u_{max}} Pr^{-2/3} \quad (10)$$

The friction factor f is used for the hydraulic characterization of the heat exchanger model. The Fanning friction factor is calculated using the Kay and London definition [12] for a fluid with constant properties, where A_c is the minimum flow area. σ is the contraction area ratio, Δp is the pressure drop and G_{min} is the mass flux in the minimal flow area, equation 11:

$$f = \frac{A_c}{A_0} \frac{\rho_m}{\rho_n} \left(\frac{2 \rho_n \Delta p}{G_{min}} - (1 + \sigma^2) \left(\frac{\rho_n}{\rho_{out}} - 1 \right) \right) \quad (11)$$

For evaluating the fin efficiency, the numerical method from Ameel *et al.* [7] can be implemented. Since Wang *et al.* [1] was using the Schmidt's method it cannot be used for finding correlations but only for validation of the fin efficiency as function of inlet velocity and fin pitch. Fourteen planes constant in the x-axis were created through the channel with 1mm distance to each other starting at the first edge ($x_1 = 0,01m$) until the end of the fin. Within each one of the planes an average temperature of the fluid phase is calculated, that will be called "bulk temperature", equation 12.

$$T_{bulk(x_i)} = \frac{\iint T(x_i, y, z) u(x_i, y, z) dy dz}{\iint u(x_i, y, z)} \quad (12)$$

For each inlet velocity, the values are obtained and a regression is performed. With this temperature trend, the local heat transfer coefficients equation 15 are calculated with the local heat flux and fin temperatures.

Equation 13:

$$h_{local}(x, y, z) = \frac{\dot{Q}_{local}(x, y, z)}{T_m(x, y, z) - T_{bulk}(x)} \quad (13)$$

An ideal fictitious heat flux is then determined as if for isothermal conditions in equation 14:

$$\dot{Q}_{iso}(x, y, z) = \iint h_{local}(x, y, z)(T_{wall} - T_{bulk}(x))dA \quad (14)$$

The fin efficiency is obtained in equation 15 by the ratio of the real to the isothermal heat flux, equation 15.

$$\eta_{fin} = \frac{\dot{Q}_{real}(x, y, z)}{\dot{Q}_{iso}(x, y, z)} \quad (15)$$

For a fin pitch of $Fp = 1,81\text{mm}$ the values of the heat transfer coefficient and pressure drop are used to calculate the Colburn and the friction factor. These values are represented in the figure 3 and compared to the results obtained by Wang *et al* [1]. The uncertainties in the reference are expressed for j and f . The figure shows a good match between experimental and simulated data obtained in this work and validates the numerical model. The independence of mesh and resolution is given through figure 4 where data for Colburn and friction factor are plotted for a coarse mesh with a total number of $3,5 \times 10^5$ volumes and in the second, finer one with $5,4 \times 10^5$ volumes.

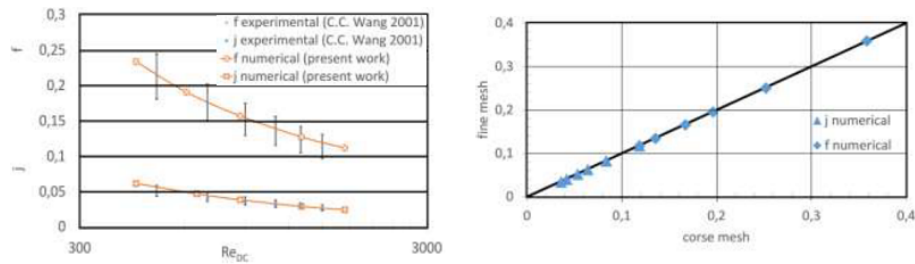


Fig. 3

Colburn and friction factor for experimental and numerical data comparison (left). Colburn and friction factor to compare results of different meshed domains (right)

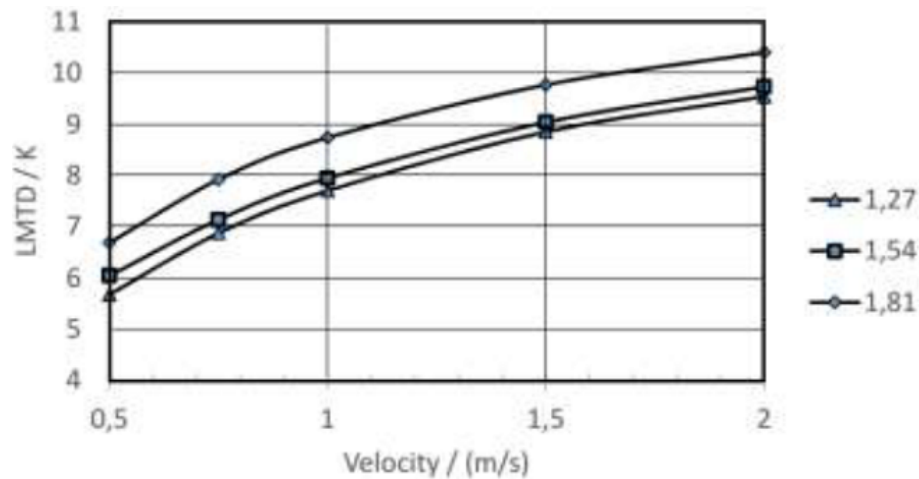


Fig. 4

Logarithmic temperature difference between in- and outlet for three different fin pitches as a function of the inlet velocity

Results and Discussion

Simulated velocities range from 0,1 to 2,0 m/s but for this article only the ones higher than 0,5 m/s are considered since only Reynolds numbers significantly larger than 200 are customary for compact heat exchangers. The numerically computed fine efficiency has a peak at around $u_{in} = 0,3$ m/s and falls rapidly for smaller velocities what could be due to the heat reversal problem occurring at the end of the fin and in the wake area. A similar problem of heat transfer reversal was studied in the work of Fiebig *et al* [13].

The simulation shows for the logarithmic temperature difference (LMTD), mass flow and pressure drop nearly uniform and predictable results. With higher fin pitch the logarithmic temperature difference in figure 5 increases as well as with rising inlet velocity of the air.

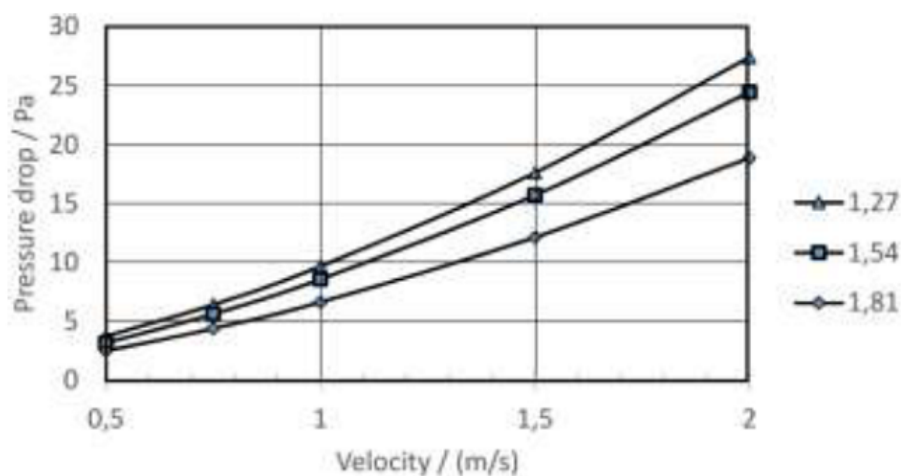


Fig. 5

Pressure drop Δp for three different fin pitches as a function of the inlet velocity

The pressure drop in figure 5 follows quadratic behavior and gets smaller for higher fin pitch. The temperature at the outlet varies between 286K and 293K. With higher inlet velocity and greater fin pitch the outlet temperature rises because there is less time for the heat transfer from fluid to solid, higher mass flow and greater ratio of channel volume to heat conduction surface.

The numerical calculated fin efficiency equation 17 is be plotted in figure 6 for every fin pitch and inlet velocity. Note that the behavior for $F_p=1,81\text{mm}$ and $F_p=1,54\text{mm}$ is quite similar while the curve for $F_p=1,27\text{mm}$ has greater values than the others except for the first data point do.

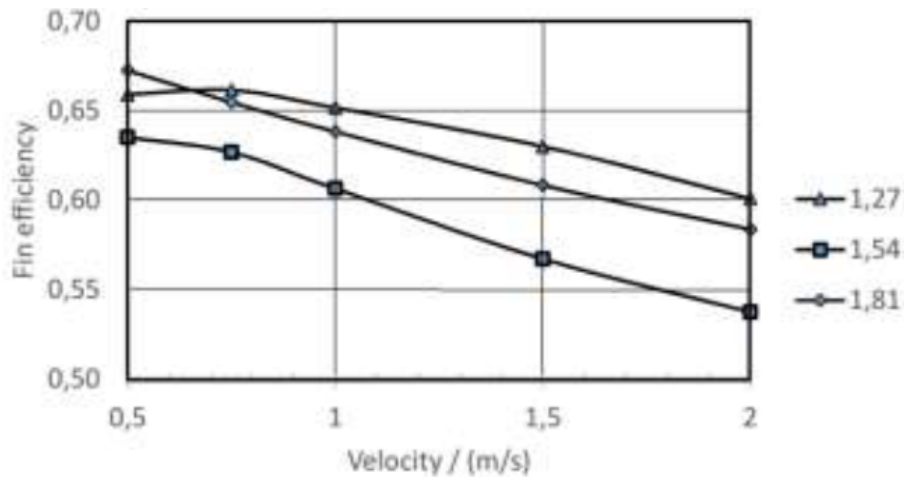


Fig. 6

The numerical calculated fin efficiency η_{fin} for different fin pitches and velocities

Wang *et al*[2] described the increase of heat transfer performance with a decrease of the fin pitch. This could be one's expectation because the ratio of real to the ideal fictitious heat flux equation 15 is getting steadily smaller. With this interpretation in mind, the heat transfer coefficient h_0 which is calculated with the equation 16 would increase. The data shows a different, not expected performance for both fin efficiency and heat transfer coefficient, figure 7.

$$h_0 = \frac{Q_{total}}{\Delta T_{log} (A_{fin} \eta_{fin} + A_{tube})} \quad (16)$$

$$J = j Re_{D_c} \quad (17)$$

$$F = f Re_{D_c}^3 \quad (17 - \text{Cont.})$$

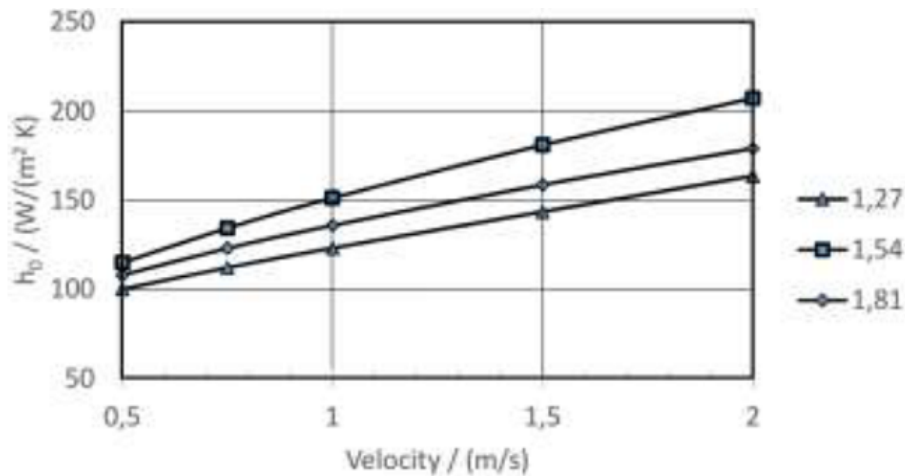


Fig. 7

Calculated heat transfer coefficient for three fin pitches for a inlet velocity of 2,0m/s

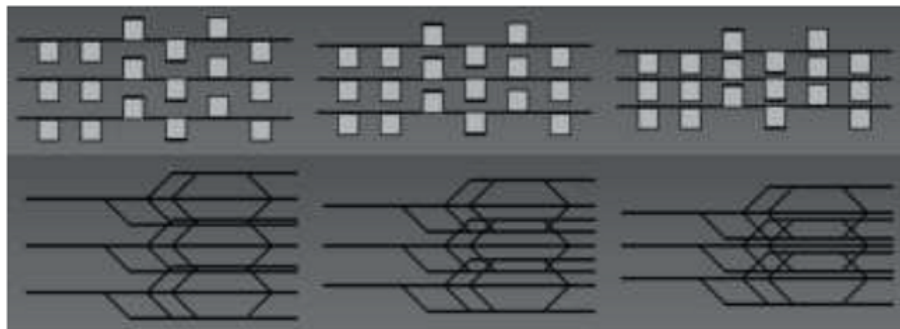


Fig. 8

Side views of the stacked fins in z -axis (upper drawing) and in x -axis (lower drawing) to see the distribution of the slits starting with highest fin pitch on the left

The quite different performance is explained due to the distribution of the strips that is shown in figure 8 where two periods of the stacks geometry are illustrated. While the one with $f_p = 1,54$ mm has a quite evenly spread geometry, the other ones have whether offset strips that are very close to the fin in z -direction or that are overlapping when looking in x -direction.

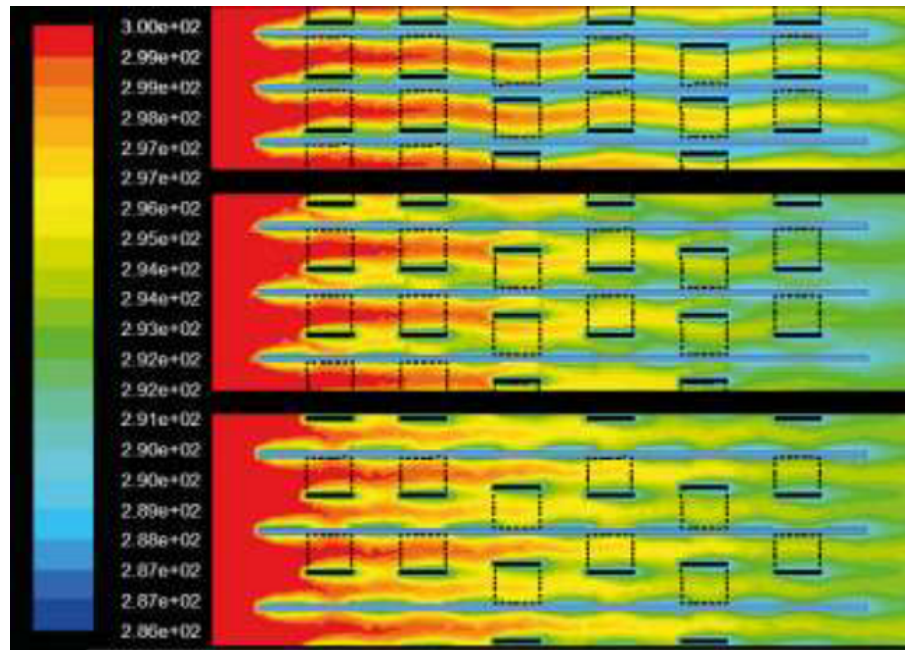


Fig. 9

Temperature field for an inlet velocity of 2,0 m/s starting with the lowest fin pitch at the top

Of course, this affects the boundary layers produced by the airflow. In figure 9 the fluids temperature is shown for all the fin pitches at an inlet velocity of 2.0m/s. For the sack of simplicity, the plate fin is drawn as a light blue rectangle without the slits and the dotted lines significance the backside of the offset strips. The upper drawing represents the temperature boundary layersfor $f_p = 1,27\text{mm}$. The strips are located very close to the eachopposite fin surface where boundary layers are already startingto develop. There the temperature gradient is lower than in thefree stream in the middle between the two fins where it is ableto flow with very few interruptions.

The configuration with $f_p = 1,81\text{mm}$ shows that the stripsfrom the both subsequent fins are one after another nearly inthe same line of stream flow. This is already better than the firstconfiguration because they are found right in theuntil then notdisturbed flow and produce new boundary layers.

Still the heat is not being exchanged very well in comparison to the second configuration where the strips lie in themiddle between the fins and alternate in their distance to them. They divide the flow various times what results in more turbulence, higher local heat fluxes and explains the special characteristics in figure 9. This configuration is also producing an about10% greater pressure loss in flow compared to the setup with $f_p = 1,81\text{mm}$ and hence to an elevated power consumption of the fan.

For the heat flux on the fin surface in figure 10, the leading edges always show the highest values and one can note, that the strips inthe middle of the fins appear to have the most noticeable change for different fin pitches. The figure may lead to the conclusionthat for $f_p = 1,54\text{mm}$ the most energy is transmitted but stillthere needs to be considered, that the highest total heat transferis monitored for the maximal fin pitch

because the mass flow as well as the temperature difference are greater at a constant inlet velocity. In addition, the tube wall exchanges energy and it gets larger with increasing fin pitch.

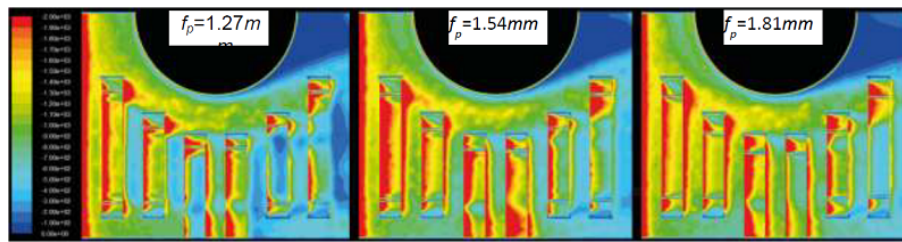


Fig. 10

Surface heat flux in W/m^2 at a velocity of 2.0 m/s starting with lowest fin pitch at the left

For the Colburn factor figure 11 the model with $f_p=1.54\text{ mm}$ shows the highest thermal performance, while for the friction factor has the worst hydraulic performance. This fits to the former observations made in this paper. We suppose an optimal fin pitch could be found for this geometry and under the conditions tested in this paper. A commonly used criterion to evaluate the thermal-hydraulic performance of a heat exchanger is the area goodness factor. It is defined as the ratio of the Colburn to the friction factor. High values of j/f are preferred as this means less frontal area for a fixed heat transfer and pressure drop according to Stone [14]. The area goodness factor is plotted in figure 12 and may suggest $f_p=1.54\text{ mm}$ as the best configuration. LaHaye *et al* [15] suggested evaluating the thermal-hydraulic performance of heat exchangers by plotting the heat transfer performance factor J as function of the pumping power factor F which is done in figure 12. The model with $f_p=1.54\text{ mm}$ has a better thermal performance for the same pumping power when compared to the other ones.

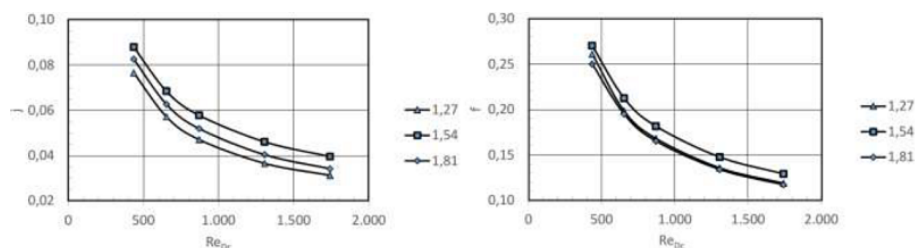


Fig. 11

Colburn factor and Friction factor for three different fin pitches as a function of the inlet velocity

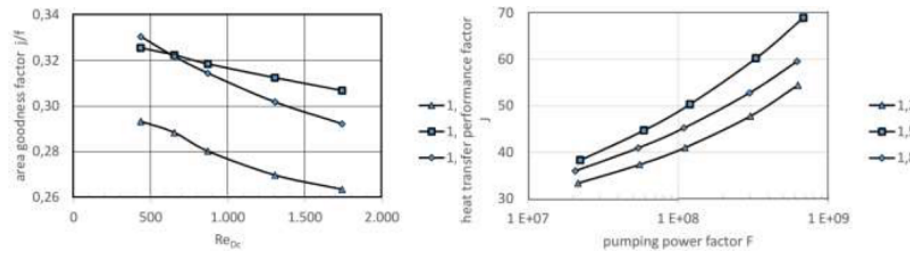


Fig. 12

Area goodness factor of the Reynolds number for different fin pitches (left) and evaluation of the thermal-hydraulic performance according to LaHaye et al. [11] (right)

Conclusions

The average heat transfer coefficient as a function of the fin pitch is at its best when the strips of one fin have the same distance to the ones of the opposite arranged fin as to their fin plate. So the highest heat flux and turbulence is achieved. The model with $f_p = 1,54$ mm has the best thermal performance for the same pumping power. The fin efficiency as a function of the inlet velocity has a peak for each fin pitch. For velocities smaller than the peak, it falls rapidly wherefore heat transfer reversal could be responsible for this fact. The calculation of fin efficiency via the Schmidt method is not adequate on interrupted fins.

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