



Journal of Aerospace Technology and
Management

ISSN: 1984-9648

secretary@jatm.com.br

Instituto de Aeronáutica e Espaço
Brasil

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Journal of Aerospace Technology and Management, vol. 9, núm. 3, julio-septiembre,
2017, pp. 389-396

Instituto de Aeronáutica e Espaço
São Paulo, Brasil

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Orthotropic Elastic Properties Assessment of Sandwich Laminates

Paulo Pedro Kenedi^{1,2}, Lucas Lisboa Vignoli³, Brenno Tavares Duarte², Fernando Cesar de Abreu Matos¹, Humberto Oberosler Terço Dias¹

ABSTRACT: Sandwich structures are nowadays widely used due its excellent characteristics of matching light weight with high bend stiffness. However the complex nature of its composition turns the estimation of its mechanical elastic properties challenging. The objective of the present article is to propose an analytical/experimental model to estimate the orthotropic elastic properties of this composite material. To accomplish this objective a theory of anisotropic plates is used together to an extensometric experimental approach. A sandwich laminate specimen, instrumented with strain gages, was submitted through pure bending loading by a mechanical testing machine on a 4-point bending apparatus. The result of this research was the determination of 4 orthotropic elastic properties of the sandwich laminate specimen, namely: the longitudinal and transversal Young moduli, the shear modulus, and the Poisson's ratio.

KEYWORDS: Sandwich laminates, Orthotropic properties, Analytical model.

INTRODUCTION

Composites materials, as sandwich laminates, have been widely used for structural purposes. These structures are anisotropic, combining external fibers (with high strength) and an internal core (with low stiffness). This combination of materials contributes to turn its stress analysis not as straightforward.

Much effort has been done to use sandwich laminates as an engineering material. Authors as Bledzki *et al.* (1999) identified elastic properties of unidirectional glass/epoxy laminates through the measuring of its natural frequencies in experimental vibration experiments. Lamboul *et al.* (2013) investigated the potential of ultrasonic Lamb waves to estimate the initial elastic modulus of glass fiber woven composite materials. The elastic modulus of the undamaged materials was estimated using time of flights (TOF) of extensional waves (also known as Lamb S0 mode) generated by a piezoelectric material. Rikards *et al.* (2001) and Rikards *et al.* (1999) used mixed numerical/experimental techniques to estimate mechanical properties of composites laminates. In the experimental part it was used real-time holography, and in the numerical part it was used the finite element method. Lecompte *et al.* (2007) have also used a mixed numerical-experimental method to identify 4 in-plane orthotropic engineering constants of composite plate materials. It was used a biaxial tensile test, performed on a cruciform test specimen. A heterogeneous displacement field was observed by a CCD camera and measured by a digital image correlation (DIC) technique. The measured displacement field and the subsequently computed strain field were compared with a finite element simulation of the same experiment. An optimization technique was then used to minimize the errors

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Received: Aug., 30, 2016 | **Accepted:** Nov., 10, 2016

between experimental and numerical approaches resulting in the estimation of 4 independent apparent engineering constants: E_x , E_y , G_{xy} , and ν_{xy} .

Other authors as Tamilarasan *et al.* (2015) evaluated mechanical properties of sandwich laminates made up of carbon fiber with epoxy resin, used as a matrix, and aluminum (AA6061-T6), used as sandwich plate, through the realization of multiple tests as: uniaxial tensile, flexural, and impact. The tests results were used by the authors to describe, in a macroscopic perspective, the different material behavior associated to different levels of strengths (axial, flexural, and impact), including the failure mechanisms.

The research of Hu and Karki (2015) involved a 2-step modeling approach to estimate the mechanical properties of 3-D carbon fiber-reinforced epoxy composites. In step 1, the microstructures of the composites were represented by 3-D microscale cylindrical, square, or hexagonal prismatic representative volume element (RVE) and were subject of multiple tests as: uniaxial tensile, lateral expansion, and transverse shear. In step 2, the overall composites were represented by a 3-D macroscale unit cell that was discretized into elements. Each element was assigned with appropriate properties, taking into account the carbon fiber orientation and volume fraction. The averaged mechanical properties of the RVEs from step 1 were input into each element accordingly. The effective orthotropic mechanical properties of the composites were predicted by uniaxial tensile tests and shear tests on the unit cell.

To improve the understanding of mechanical behavior of sandwich laminates, this article proposes a combined model, using analytical and experimental approaches, to estimate 4 orthotropic mechanical elastic properties of sandwich laminates. The analytical part uses a theory of anisotropic plates and the experimental one, strain gages to obtain the strains generated on a sandwich laminate specimen, submitted to pure bending on a 4-point bending test. Note that although, in technical literature, it is more common the utilization of unidirectional composite materials, in this article a woven (fabric) sandwich laminate was used to obtain the longitudinal and the transversal Young moduli, the shear modulus, and the Poisson's ratio mechanical properties. In the next section the analytical models are described.

ANALYTICAL MODELS

Some initial information, as the 4-point bending apparatus and the specimen principal dimensions are provided before the analytical models description are given. Figure 1 shows a

schematic representation of the sandwich laminate specimen, instrumented with strain gages, submitted to pure bending in its central region (between upper roller anvils), on a 4-point bending apparatus.

Note in Fig. 1a that the dark blue color represents the external polymeric matrix, with high strength carbon fibers ($0^\circ/90^\circ$)/epoxy (with 0.6 mm of thickness each) and the light blue color represents the low stiffness PVC foam (with 3.8 mm of thickness). The distance between the upper roller anvil and the immediate adjacent lower roller anvil is $l = 50$ mm and the distance between the lower roller anvils is $L = 170$ mm.

Figure 2 shows the sandwich laminate specimen schematic drawing, with specimen nominal dimensions and the used coordinate systems. θ is an angle, in counterclockwise direction, between local and global systems, beginning with $\theta = 0^\circ$ on global axis direction 11.

Two analytical models are provided to estimate orthotropic elastic constants of composite sandwich specimens, both considering Euler-Bernoulli hypothesis. The first model is a simplified approach which disregards the core material stiffness and the second one is a complete analytical solution that considers

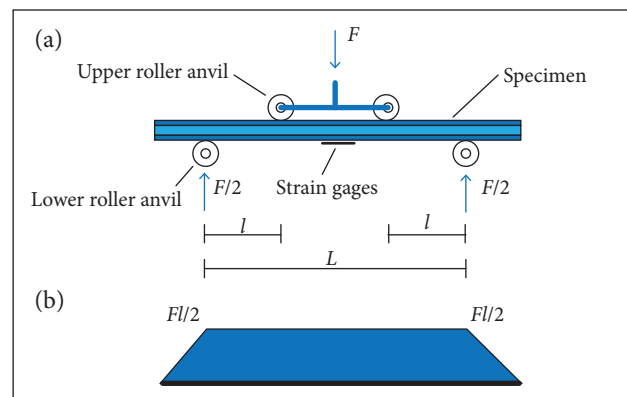


Figure 1. (a) Schematic representation of the 4-point bending apparatus with specimen (lateral view); (b) Correspondent bending moment diagram.

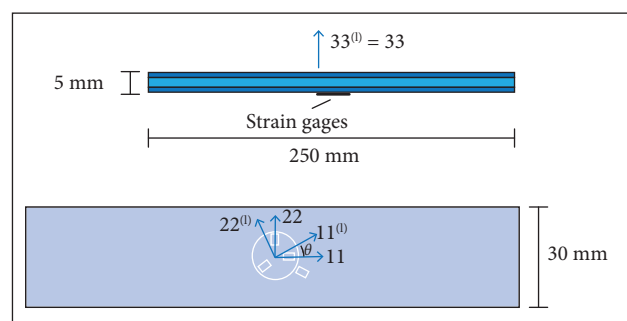


Figure 2. Composite specimen schematic drawing at a third angle projection (unscaled).

the stiffness of every layer, including the different elastic moduli, in tension and in compression, of the core material. Both models depend on the experimental strain acquisition. Regardless of the layers anisotropy, it is assumed that the sections remain plane and perpendicular to the neutral axis, being enough to describe the normal strain variation along the cross section. A constitutive relation and the equilibrium requirements are additionally provided to implement the stress calculation, as in Lekhnitskii (1981).

SIMPLIFIED ANALYTICAL MODEL

The simplified model is divided in stress and strain parts. First, the stress is estimated by application of a well-known mechanics of solids equation and the proposal of a constitutive relationship. Then, with the input of experimental strains, 4 elastic orthotropic constants are estimated.

At the central region of a 4-point bending setup, between upper roller anvils, the shear force vanishes and it is considered that this region is only submitted to pure bending loading. Hence, the unique non-null global stress component on the bottom of the beam in this region is, as in Crandall *et al.* (1978),

$$\sigma_{11} = \frac{M_{22}}{I_{22}} x_3 \quad (1)$$

where: x_3 is the vertical distance between the neutral axis and a given point of interest; M_{22} is the pure moment around 22 axis (see Fig. 2); I_{22} is the cross section second moment of area relative to 22 axis:

$$M_{22} = Fl/2$$

$$I_{22} = \frac{b}{12} \left[(t_c + 2t_l)^3 - (t_c)^3 \right] \quad (2)$$

$$x_3 = t_l + t_c/2$$

where: F is the applied load; l is the distance between the upper roller anvil and the immediate adjacent lower roller anvil (see Fig. 1a); t_l is the woven lamina thickness; t_c is the core thickness; b is the specimen width.

Note that for this simplified analytical model it is supposed that all loads are only resisted by the woven laminas, since the core material has a much lower stiffness (Kassapoglou 2010). This simplified approach turns the σ_{11} estimation straightforward.

The normal local strains, at a general given angle θ , in relation to the global axes, shown in Fig. 2, can be estimated as in Lekhnitskii (1981):

$$\varepsilon_{11}^{(l)} = C_{11}^{(l)} \sigma_{11}^{(l)} + C_{12}^{(l)} \sigma_{22}^{(l)} + C_{16}^{(l)} \sigma_{12}^{(l)} \quad (3)$$

where the local stress components $\sigma_{11}^{(l)}$, $\sigma_{22}^{(l)}$, and $\sigma_{12}^{(l)}$ can be obtained through the use of Mohr's circle as in Crandall *et al.* (1978):

$$\sigma_{11}^{(l)} = \frac{\sigma_{11}}{2} (1 + \cos 2\theta)$$

$$\sigma_{22}^{(l)} = \frac{\sigma_{11}}{2} (1 - \cos 2\theta) \quad (4)$$

$$\sigma_{12}^{(l)} = -\frac{\sigma_{11}}{2} \sin 2\theta$$

and the local compliance components can be estimated as in Staab (2015):

$$C_{11}^{(l)} = C_{11} \cos^4 \theta + (2C_{12} + C_{66}) \sin^2 \theta \cos^2 \theta + C_{22} \sin^4 \theta \quad (5)$$

$$C_{12}^{(l)} = C_{12} + (C_{11} + C_{22} - 2C_{12} - C_{66}) \sin^2 \theta \cos^2 \theta \quad (6)$$

$$C_{16}^{(l)} = (2C_{11} - 2C_{12} - C_{66}) \sin \theta \cos^3 \theta - (2C_{22} - 2C_{12} - C_{66}) \sin^3 \theta \cos \theta \quad (7)$$

Considering that the material coordinate system is coincident with the global coordinate axes, the reduced global compliance matrix C , also considering just the in-plane components and its orthotropic symmetry for a woven lamina, can be written as in Staab (2015):

$$C = \begin{bmatrix} C_{11} & C_{12} & C_{16} \\ C_{12} & C_{22} & C_{26} \\ C_{16} & C_{26} & C_{66} \end{bmatrix} = \begin{bmatrix} 1/E_1 & -\nu_{12}/E_1 & 0 \\ -\nu_{12}/E_1 & 1/E_2 & 0 \\ 0 & 0 & 1/G_{12} \end{bmatrix} \quad (8)$$

where: E_1 and E_2 are the lamina elastic moduli for, respectively, 11 and 22 axis global directions; G_{12} is the lamina shear modulus; ν_{12} is the lamina Poisson's ratio.

Using the compliance matrix components, shown in Eq. 8, Eqs. 5 – 7 can be rewritten for plane stress case, as in Staab (2015):

$$C_{11}^{(l)} = \frac{1}{E_1} \cos^4 \theta + \left(-2 \frac{\nu_{12}}{E_1} + \frac{1}{G_{12}} \right) \sin^2 \theta \cos^2 \theta + \frac{1}{E_2} \sin^4 \theta \quad (9)$$

$$C_{12}^{(l)} = -\frac{\nu_{12}}{E_1} + \left(\frac{1 + 2\nu_{12}}{E_1} + \frac{1}{E_2} - \frac{1}{G_{12}} \right) \sin^2 \theta \cos^2 \theta \quad (10)$$

$$C_{16}^{(l)} = 2 \left(\frac{1 + \nu_{12}}{E_1} - \frac{1}{2G_{12}} \right) \sin \theta \cos^3 \theta - 2 \left(\frac{1}{E_2} + \frac{\nu_{12}}{E_1} - \frac{1}{2G_{12}} \right) \sin^3 \theta \cos \theta \quad (11)$$

Notice that Eq. 3 can be rewritten in a matricial form, with 4 lines, one per strain gage of angle θ_n .

$$\begin{bmatrix} \varepsilon_{11}^{(l)}|_{\theta_1} \\ \varepsilon_{11}^{(l)}|_{\theta_2} \\ \varepsilon_{11}^{(l)}|_{\theta_3} \\ \varepsilon_{11}^{(l)}|_{\theta_4} \end{bmatrix} = \begin{bmatrix} f_1(\theta_1) & f_2(\theta_1) & f_3(\theta_1) & f_4(\theta_1) \\ f_1(\theta_2) & f_2(\theta_2) & f_3(\theta_2) & f_4(\theta_2) \\ f_1(\theta_3) & f_2(\theta_3) & f_3(\theta_3) & f_4(\theta_3) \\ f_1(\theta_4) & f_2(\theta_4) & f_3(\theta_4) & f_4(\theta_4) \end{bmatrix} \begin{bmatrix} 1/E_1 \\ \nu_{12}/E_1 \\ 1/E_2 \\ 1/G_{12} \end{bmatrix} \frac{\sigma_{11}}{2} \quad (12)$$

where:

$$f_1(\theta) = 2 \cos^2 \theta (2 \cos^2 \theta - 1)^2 \quad (13)$$

$$f_2(\theta) = 2 \cos^3 2\theta - \cos 2\theta - 1$$

$$f_3(\theta) = -8 \sin^4 \theta (\sin^2 \theta - 1) \quad (14)$$

$$f_4(\theta) = \cos 2\theta \sin^2 2\theta$$

The final equation can be obtained rearranging Eq. 12, substituting the result of Eq. 1 and the strain-gages angles (from θ_1 to θ_4), and its respective experimental values (from $\varepsilon_{11}^{(l)}|_{\theta_1}$ to $\varepsilon_{11}^{(l)}|_{\theta_4}$).

$$\begin{bmatrix} 1/E_1 \\ \nu_{12}/E_1 \\ 1/E_2 \\ 1/G_{12} \end{bmatrix} = \begin{bmatrix} f_1(\theta_1) & f_2(\theta_1) & f_3(\theta_1) & f_4(\theta_1) \\ f_1(\theta_2) & f_2(\theta_2) & f_3(\theta_2) & f_4(\theta_2) \\ f_1(\theta_3) & f_2(\theta_3) & f_3(\theta_3) & f_4(\theta_3) \\ f_1(\theta_4) & f_2(\theta_4) & f_3(\theta_4) & f_4(\theta_4) \end{bmatrix}^{-1} \begin{bmatrix} \varepsilon_{11}^{(l)}|_{\theta_1} \\ \varepsilon_{11}^{(l)}|_{\theta_2} \\ \varepsilon_{11}^{(l)}|_{\theta_3} \\ \varepsilon_{11}^{(l)}|_{\theta_4} \end{bmatrix} \frac{2}{\sigma_{11}} \quad (15)$$

With Eq. 15 it is possible to estimate 4 elastic orthotropic constants (E_1 , E_2 , ν_{12} , and G_{12}). A flowchart is shown in Fig. 3 as a form to review the steps necessary to apply the simplified analytical model.

The flowchart of Fig. 3 shows the straightforward utilization of the simplified analytical model. In fact to estimate the 4 orthotropic elastic constants E_1^{SAM} , E_2^{SAM} , ν_{12}^{SAM} , G_{12}^{SAM} the angles of the strain gages have to be entered in relation to the specimen longitudinal axis $\theta_1, \theta_2, \theta_3$, and θ_4 ; the correspondent experimental strain values $\varepsilon_{11}^{(l)}|_{\theta_1}$, $\varepsilon_{11}^{(l)}|_{\theta_2}$, $\varepsilon_{11}^{(l)}|_{\theta_3}$, and $\varepsilon_{11}^{(l)}|_{\theta_4}$; and the estimative of longitudinal bottom specimen stress σ_{11} into the linear system of Eq. 15 to obtain the estimative E_1^{SAM} , E_2^{SAM} , ν_{12}^{SAM} , G_{12}^{SAM} . The index "SAM" is the abbreviation of simplified analytical model. The complete analytical model is provided in the next subsection.

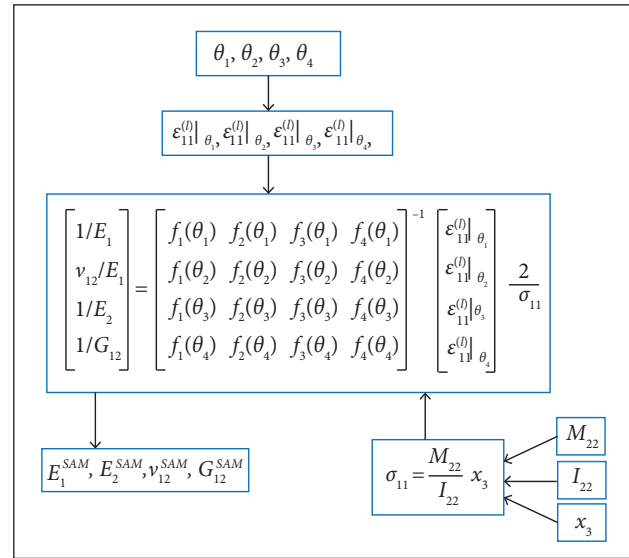


Figure 3. Simplified analytical model flowchart.

COMPLETE ANALYTICAL MODEL

In this section the influence of the core on the estimative of the elastic orthotropic constants is carefully studied. To accomplish this objective a complete analytical model is presented considering the core asymmetric tension/compression behavior, as in Kenedi *et al* (2016). The procedure presented next is based on solid mechanics theory and further discussions can be found in Crandall *et al.* (1978) as well as in Vasiliev and Morozov (2001).

In Fig. 4 $E_c^{(t)}$ and $E_c^{(c)}$ are, respectively, the core elastic modulus in tension and in compression. As the core tension properties are different from the compression ones, the stress distribution is no more symmetric, turning the neutral axis position \bar{x}_3 (see Fig. 4) an additional parameter to be taken

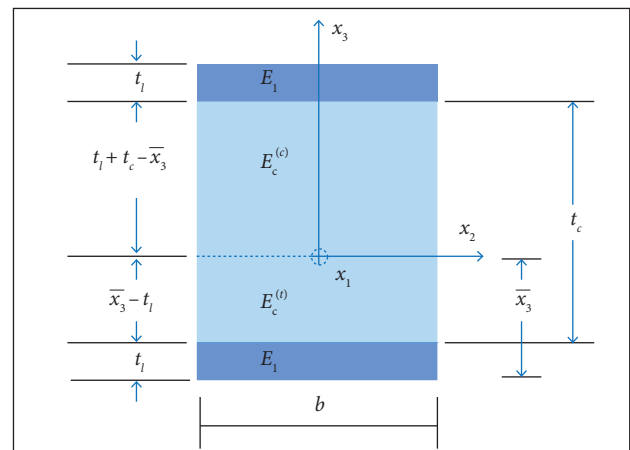


Figure 4. Sandwich laminates specimen transversal cross-section dimensions.

into account by imposing the equilibrium conditions along the beam cross section. In fact $E_c^{(t)}$ and $E_c^{(c)}$ must be used as 2 limiting conditions with minimum and maximum properties, called, respectively, in this article, as lower and upper core bounds, which are available in Table 1. These lower and upper properties cover the range of elastic moduli available in technical literature for the core material, for instance, PVC foam. Note that to use both models there is no need for having any estimative of the lamina elastic moduli E_l , which can be considered an advantage due to the great variability of this property in function of parameters like: fraction volume, epoxy type, fabric patterns and fabrication process.

Table 1. Range of elastic modulus core material.

	Lower bound core (lo)	Upper bound core (up)
$E_c^{(t)}$ (MPa)	49	320
$E_c^{(c)}$ (MPa)	40	400

Source: Diab Group and Divinycell (2016).

Taking the asymmetric characteristic into account, considering that the section remains plane and perpendicular to the neutral axis and using the linear elastic constitutive relation, the force equilibrium along the cross section can be written as (Crandall *et al.* 1978):

$$\int \sigma_{11} dA = \int E \varepsilon_{11} dA = \frac{1}{\rho} \int E x_3 dA = 0 \quad (16)$$

where: ρ is the radius of curvature, which is supposed to be constant on the cross section.

Equation 16 is the definition of the centroid in non-homogeneous materials, which can be rewritten as:

$$\begin{aligned} & (E_c^{(t)} - E_c^{(c)}) \bar{x}_3^2 + 2 \left[(2E_1 - E_c^{(t)} + E_c^{(c)}) t_l + E_c^{(c)} t_c \right] \bar{x}_3 - \\ & - \left[(4E_1 - E_c^{(t)} + E_c^{(c)}) t_l^2 + 2(E_1 + E_c^{(c)}) t_l t_c + E_c^{(c)} t_c^2 \right] = 0 \end{aligned} \quad (17)$$

Note that Eq. 17 is a second-order polynomial function but it is not completely defined once E_l is still unknown. $E_c^{(t)}$ and $E_c^{(c)}$ are first substituted by lower bounds and then by upper ones to verify the solution range. Since the neutral axis expression is obtained, the moment equilibrium on the cross section, following the same assumptions used for the force equilibrium, presented previously, is defined as Crandall *et al.* (1978):

$$\int x_3 \sigma_{11} dA = \frac{1}{\rho} \int E x_3^2 dA = M_{22}, \quad (18)$$

defining equivalent flexural rigidity as:

$$EI_{eq} \left(E_c^{(c)}, E_c^{(t)}, E_1 \right) = \int E x_3^2 dA \quad (19)$$

This quantity can be computed solving the integral or alternatively using the parallel axis theorem as

$$\begin{aligned} EI_{eq} \left(E_c^{(c)}, E_c^{(t)}, E_1 \right) &= E_1 I_{22}^{(l,t)} + E_c^{(t)} I_{22}^{(f,t)} + \\ &+ E_c^{(c)} I_{22}^{(f,c)} + E_1 I_{22}^{(l,c)} \end{aligned} \quad (20)$$

where:

$$I_{22}^{(l,t)} = \frac{b t_l^3}{12} + (b t_l) \left(\bar{x}_3 - \frac{t_l}{2} \right)^2 \quad (21)$$

$$\begin{aligned} I_{22}^{(f,c)} &= \frac{b t_l^3}{12} + (b t_l) \left(\frac{3t_l}{2} + t_c - \bar{x}_3 \right)^2 \\ I_{22}^{(f,t)} &= \frac{b}{12} (\bar{x}_3 - t_l)^3 + \left[b (\bar{x}_3 - t_l) \right] \left(\frac{\bar{x}_3 - t_l}{2} \right)^2 \\ I_{22}^{(f,c)} &= \frac{b}{12} (t_l + t_c - \bar{x}_3)^3 + \left[b (t_l + t_c - \bar{x}_3) \right] \left(\frac{t_l + t_c - \bar{x}_3}{2} \right)^2 \end{aligned} \quad (22)$$

The upper indexes of the second moment of area are related to: the first index for the material l represents the lamina and f represents the foam (at core); the second index for the stress state t represents tension and c represents compression. Note that to maintain compact the notation in Eqs. 21 and 22, $\bar{x}_3 (E_c^{(c)}, E_c^{(t)}, E_1)$ and $I_{22}^{(c)} (E_c^{(c)}, E_c^{(t)}, E_1)$ are written, respectively, as \bar{x}_3 and $I_{22}^{(c)}$.

Thus, the radius of curvature expression is obtained from direct substitution in Eq. 18. Hence, the normal stress on the laminas of the sandwich laminate is calculated as

$$\sigma_{11} \left(E_c^{(c)}, E_c^{(t)}, E_1 \right) = E_1 \frac{M_{22}}{EI_{eq} \left(E_c^{(c)}, E_c^{(t)}, E_1 \right)} \bar{x}_3 \left(E_c^{(c)}, E_c^{(t)}, E_1 \right) \quad (23)$$

In Fig. 5 a flowchart of the complete analytical model is presented and shows a little more laborious implementation than the flowchart of Fig. 3. Note that this flowchart needs the results of the former one (Fig. 3), E_1^{SAM} , E_2^{SAM} , ν_{12}^{SAM} , G_{12}^{SAM} , as an initial trial to solve a linear system through the use of the “find” command of Mathcad software to calculate E_l and,

consequently, E_2 , ν_{12} , and G_{12} . In fact, the complete analytical model has to be applied twice: first the lower bond core properties are used (with an index $_{lo}$), shown in Table 1, for the $E_c^{(t)}$ and $E_c^{(c)}$ properties to estimate: \bar{x}_3 and $EI_{eq}(E_c^{(c)}, E_c^{(t)}, E_1)$, which are substituted in $\sigma_{11}(E_c^{(c)}, E_c^{(t)}, E_1)$. Inputting the angles of the strain gages θ_1 , θ_2 , θ_3 , and θ_4 and the correspondent experimental strain values $\epsilon_{11}^{(l)}|_{\theta_1}$, $\epsilon_{11}^{(l)}|_{\theta_2}$, $\epsilon_{11}^{(l)}|_{\theta_3}$, $\epsilon_{11}^{(l)}|_{\theta_4}$ into Eq. 15 and activating the “find” command of Mathcad software, the $E_{1_{lo}}^{CAM}$, $E_{2_{lo}}^{CAM}$, $\nu_{12_{lo}}^{CAM}$, $G_{12_{lo}}^{CAM}$ results are obtained. Note that the index “CAM” is the abbreviation of complete analytical model. At the second time the upper bond core properties are used (with an index “up”), as shown in Table 1, for the $E_c^{(t)}$ and $E_c^{(c)}$ properties to estimate: \bar{x}_3 and $EI_{eq}(E_c^{(c)}, E_c^{(t)}, E_1)$, which are substituted in $\sigma_{11}(E_c^{(c)}, E_c^{(t)}, E_1)$; the same sequence that is done in the first time is repeated, resulting in $E_{1_{up}}^{CAM}$, $E_{2_{up}}^{CAM}$, $\nu_{12_{up}}^{CAM}$, $G_{12_{up}}^{CAM}$.

In the next section the experimental approach necessary to obtain 4 strain measurements in 4 distinct angles is explained.

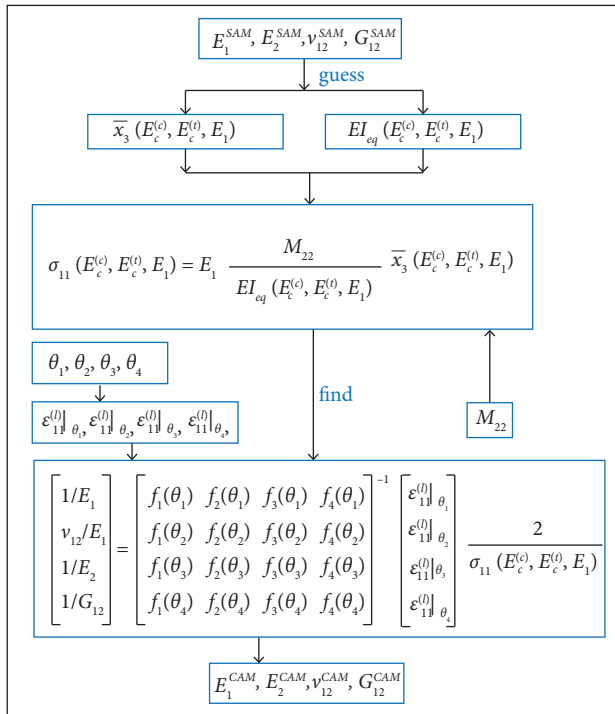


Figure 5. Complete analytical model flowchart.

EXPERIMENTAL APPROACH

In this section it is made a description of the equipment utilized in the experimental approach. All the sensors, transducers, data acquisition equipment, and materials testing machine used in the experiment to obtain the strain signals are

described in details. In Fig. 6 it is shown the 4-point bending test apparatus, used in an INSTRON material testing machine. It was used a 5-kN load cell and the LVDT of the test machine. The test was conducted monotonically under load control as stated in ASTM C393-00.

Note, in Fig. 6, the instrumented specimen is lying on a 4-point bending apparatus. The strain gages are bonded at the bottom side of the specimen, as can be seen in Fig. 6b. It is possible to see a detailed view of the strain gages specimen instrumentation in Fig. 7a, where 3 of these strain gages are provided by an extensometric rosette and the fourth strain gage is provided by a uniaxial one. In Fig. 7b the denomination and the angle of each strain gage are presented.

A Spider 8 data acquisition system was used together with a HBM Catman Software to set up the experimental strain data acquisition. It was used a Kyowa extensometric rosette KFG-3-120-D28-11, as well as a Kyowa uniaxial strain gage

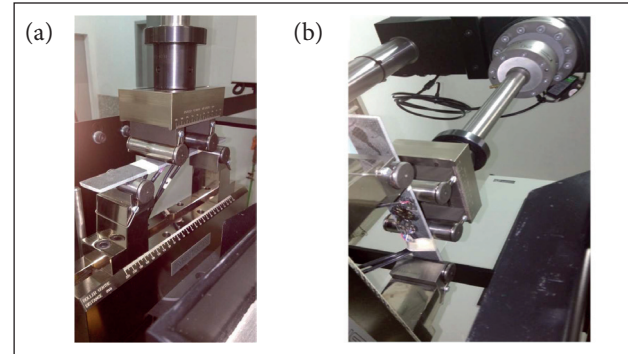


Figure 6. 4-point bending apparatus with instrumented specimen. (a) General view; (b) Specimen underside view.

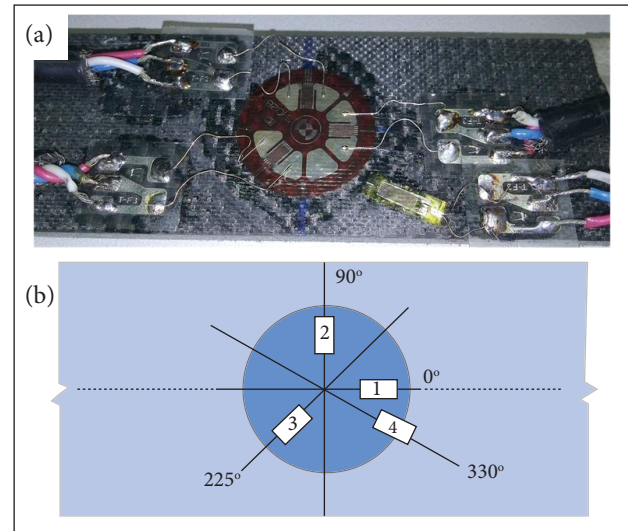


Figure 7. (a) Strain gages details; (b) Strain gage positions.

KFG-5-120-C1-11, bonded, at a sanding-prepared surface of the bottom lamina of the sandwich laminate specimen, with the utilization of a proper cyanoacrylate glue. Note in the detailed view (Fig. 7a), the utilized cabling. More information about the utilization of strain gages to measure strains on composites materials can be accessed at the fourth chapter of Staab (2015).

The objective of this test was to experimentally acquire the strain signals at 4 different angles, 0° (longitudinal and aligned to global axis 11), 90° (transversal), 225° and 330° (both inclined in relation to specimen longitudinal axis). The pure bending load was selected in order to make the strain gages positioning less critical, principally on the longitudinal direction.

Note that for analytical models, with the objective of using a more compact notation the strain gages were numbered from 1 to 4. The index equivalence between analytical and experimental notation is: 1 = 0° , 2 = 90° , 3 = 225° , and 4 = 330° , as also can be seen in Fig. 7b.

RESULTS AND DISCUSSION

In this section the experimental strain signals outputs of the 4 strain gages are provided. Figure 8 shows the strain signals outputs of the strain gages bonded at the central part of the bottom side of the specimen submitted to a pure bending on a 4-point bending test.

The data acquisition rate used in this test was selected on 50 Hz for all the 4 channels, or, in other words, 50 samples per second for each channel. Many of the acquired points of Fig. 8 were suppressed, without any loss of relevant information, to turn the experimental results clearer. It was imposed

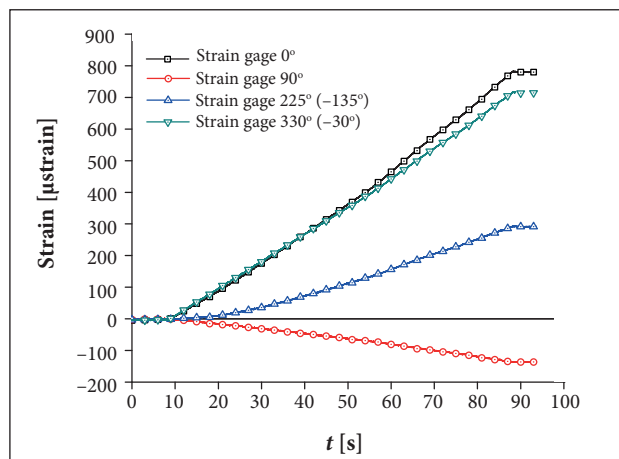


Figure 8. Strains experimental results.

a displacement rate of 1 mm/min until the load reached the force plateau of 40 N and a correspondent transversal displacement of 1.3 mm.

Note that the monotonic increasing transversal force was arbitrarily stopped at a plateau of 40 N and transversal displacement of 1.3 mm because the specimen was perceptively beginning to assume a curved configuration, which could conflict with the applicability of the mechanics of solids theory, used for instance in Eqs. 1 and 2. It would not be expected to have any perceptible geometric differences between loaded and unloaded configurations. Also it was used the high-quality load cell, already mounted on the test machine, which provided a load signal with a very low noise level, even for such low loads.

As expected the major strains occurred at 0° strain gage, aligned to the longitudinal direction (parallel to 11 global axis); at 225° and 330° the strains were at intermediary levels and, at 90° strain gage, which is the transversal one (parallel to 22 global axis), under the influence of Poisson's rate effect, it resulted in negative strains.

For this study the following strains results were used for both analytical models: $\varepsilon_{11}^{(l)}|_{\theta_1} = 780 \mu\epsilon$, $\varepsilon_{11}^{(l)}|_{\theta_2} = -136 \mu\epsilon$, $\varepsilon_{11}^{(l)}|_{\theta_3} = 292 \mu\epsilon$ and $\varepsilon_{11}^{(l)}|_{\theta_4} = 715 \mu\epsilon$.

Table 1 shows the elastic modulus range of the core material (PVC foam) obtained in technical literature (Diab Group and Divinycell 2016).

Note that as previously explained in the section "Complete Analytical Model", there is a need for the lower bond properties (with an index "lo") $E_c^{(l)}$ and $E_c^{(c)}$ and the upper bond properties (with an index "up") $E_c^{(t)}$ and $E_c^{(c)}$. These properties were provided in Table 1: $E_{c,lo}^{(t)} = 49 \text{ MPa}$, $E_{c,lo}^{(c)} = 40 \text{ MPa}$, $E_{c,up}^{(t)} = 320 \text{ MPa}$, and $E_{c,up}^{(c)} = 400 \text{ MPa}$. As already commented to use both models there is no need for having any estimative of the lamina elastic moduli E_1 , which is an advantage since the great variability of this property is a function of parameters as: fraction volume, epoxy type, fabric patterns, and fabrication process.

In Table 2 there are the results of application of the models presented in section "analytical models" to estimate the orthotropic elastic properties of the sandwich laminates specimens: E_1 , E_2 , G_{12} , and ν_{12} .

Analyzing the proposed model results written in Table 2, it is clear that the simplified analytical model produced good results if compared to the range of results of the complete analytical model (between lower and upper bounds). Nevertheless the

complete analytical model could give a better contribution to the estimation of elastic mechanical properties of sandwich laminates which have a stiffer core.

Table 2. Orthotropic elastic properties of the sandwich laminates specimens.

Model	Mechanical properties			
	E_1 [GPa]	E_2 [GPa]	ν_{12}	G_{12} [GPa]
Simplified (SAM)	18.3	19.8	0.17	5.15
Complete (CAM) (lower bound)	18.3	19.8	0.17	5.16
Complete (CAM) (upper bound)	18.6	20.2	0.17	5.25

CONCLUSIONS

Two combined analytical/experimental models were proposed to estimate 4 orthotropic mechanical elastic properties of sandwich laminates. The analytical part uses a theory of

anisotropic plates and the experimental one, strain gages to obtain the strains generated on a sandwich laminate specimen submitted to a pure bending load on a 4-point bending test. As a result, the 4 orthotropic mechanical elastic properties of sandwich laminates were estimated: the longitudinal and the transversal Young moduli, the shear modulus, and the Poisson's ratio. In short, it was proposed a relatively inexpensive and straightforward form of estimating 4 orthotropic mechanical elastic properties of sandwich laminates, which adds a little contribution to the development of this challenging area of composite materials.

AUTHOR'S CONTRIBUTION

Kenedi PP and Vignoli LL cowrote the main text. Duarte BT, Matos FCA e Dias HOT supported the experimental parts. All authors discussed the results and commented on the manuscript.

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